

"La Physique dans le Plan "

Les Houches, 5 Mars , 2010

Y-system for the spectrum of AdS/CFT

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(ENS, Paris)

with N.Gromov and P.Vieira,

with N.Gromov, A.Kozak and P.Vieira

with N.Gromov and Z.Tsuboi

with S.Leurent

arXiv:0812.5091

arXiv:0901.3753

arXiv:0906.4240

arXiv:0902.4458

arXiv:1002.3981

to appear

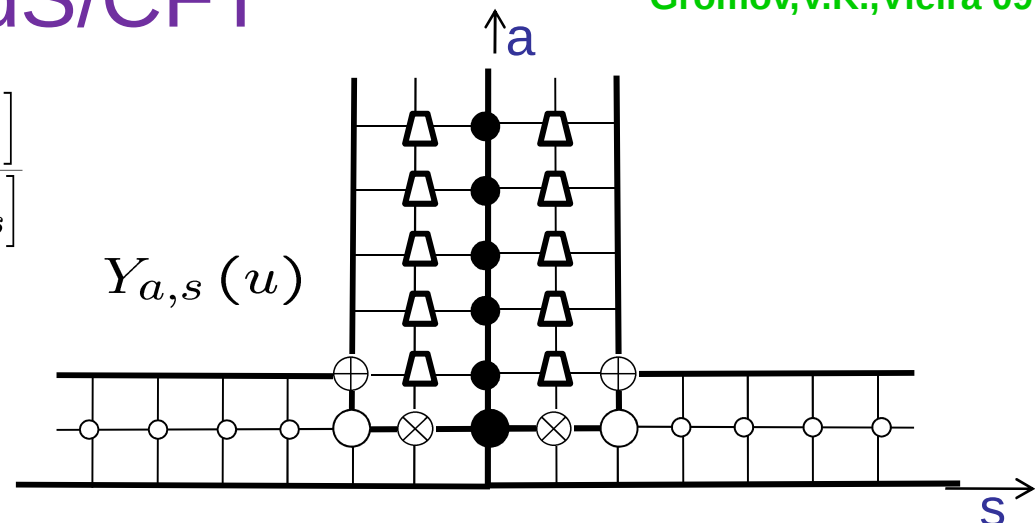
Outline

- Integrability for QFT's in dimensions $D > (1+1)$?
- Possible for some CFT's at $D=3,4$ (and may be more...)
- Based on AdS/CFT duality: $(1+1)D$ superstring sigma models on AdS-type backgrounds
- String σ -models are 2D reparametrization invariant (more than conformal): β -function = 0, $c=26$
- In light-cone gauge it looks like a “massive” non-relativistic σ -model
- All popular 2D integrability tools applicable: S-matrix, asymptotic Bethe ansatz (ABA), TBA for finite size spectrum, ...
- Y-system for AdS/CFT

Gromov, V.K., Vieira'09

$$\frac{Y_{a,s}^+ Y_{a,s}^-}{Y_{a+1,s} Y_{a-1,s}} = \frac{[1 + Y_{a,s+1}][1 + Y_{a,s-1}]}{[1 + Y_{a+1,s}][1 + Y_{a-1,s}]}$$

Calculates anomalous dimensions
of all local operators in planar $N=4$
SYM theory at any coupling



Integrability on SYM side

- Origins of YM integrability: Lipatov's reggeon Hamiltonian

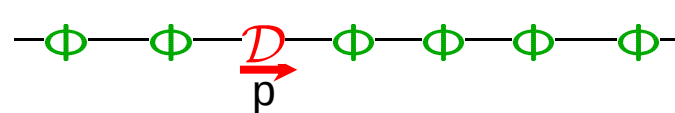
Lipatov
Faddeev, Korchemsky

- 1,2,3,4,... ∞ –loop integrability in N=4 SYM

Minahan, Zarembo,

- Exact spectrum of impurities on the SYM spin chain

Beisert, Kristijanssen,
Staudacher



$$\epsilon(p) = \sqrt{1 + \lambda \sin^2 \frac{p}{2}}$$

Fiamberti, Zieg, Zanon,
Santambogio
Beisert, Dippel, Staudacher

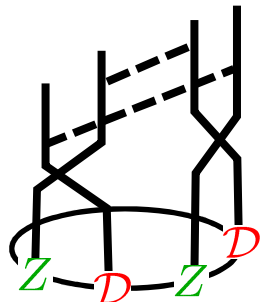
- All loop asymptotic Bethe ansatz for long operators.
PSU(2,2|4) S-matrix, Cusp dimension at all loops.

Staudacher, Beisert
Beisert, Eden, Staudacher,
Janik, Hernandez, Lopez

- Direct 4-loop calculations and wrapping for short operators

Bern, Kosover, Smirnov, Dixon
Santambrogio, Zanon,
Velizhanin

$$\Delta_{Konishi} = 2 + 12\lambda - 48\lambda^2 + 336\lambda^3 - [(2820 + 288\zeta(3)) + (324 + 864\zeta(3) - 1440\zeta(5))] \lambda^4 \\ + [(26508 + 4320\zeta(3) + 2880\zeta(5)) + (-11340 + 2592\zeta(3) - 5184\zeta(3)^2 - 11520\zeta(5) + 30240\zeta(7))] \lambda^5 + O(\lambda^6)$$



$$\mathcal{O}_{Konishi} = \text{Tr} [\mathcal{D}, Z]^2$$

... reproduced from Y-system

Gromov, V.K., Vieira
Arutyunov, Frolov
Balog, Hegedus

Integrability on string side

- Classical integrability of Green-Schwarz-Metsaev-Tseytlin string σ -model on $AdS_5 \times S^5$ and Lax pair
Berna,Roiban,Polchinski
- Finite gap solution of string σ -model and quasilcassics
V.K.,Marshakov, Minahan, Zarembo, Beisert,V.K.,Sakai,Zarembo Gromov,Vieira
- Early attempts of non-perturbative quantization
Arutyunov,Frolov,Staudacher Beisert,Staudacher
- Dressing factor
Janik, Hernandez,Lopez Beisert,Eden,Staudacher,
- Konishi dimension 4- and 5-loop calculations from string (agrees with direct 4-loop SYM perturbation theory)
Bajnok,Janik,Lukowski
- Y-system for AdS/CFT: the spectral equation for exact dimensions of planar $N=4$ SYM at any YM coupling.
Gromov , V.K, Vieira
- TBA for AdS/CFT, confirming and clarifying the Y-system
Bombardelli,Fioravanti, Tateo Gromov , V.K, Vieira Arutyunov, Frolov
- New integrable AdS/CFT's: AdS_4/CFT_3 , AdS_3/CFT_2 ,.....
Minahan,Zarembo, Zarembo

N=4 SYM and a string on $AdS_5 \times S^5$

$$\mathcal{S}_{SYM} = \frac{1}{\lambda} \int d^4x \text{Tr} \left(F^2 + (\mathcal{D}\Phi)^2 + \bar{\Psi}\mathcal{D}\Psi + \bar{\Psi}\Phi\Psi + [\Phi, \Phi]^2 \right)$$

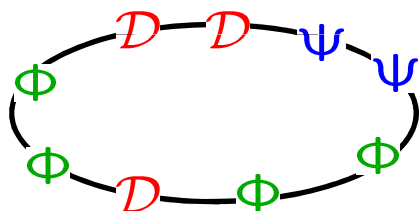
Maldacena'97
Gubser, Polyakov, Klebanov'98
Witten'98

$$\mathcal{S}_{sigma} = \sqrt{\lambda} \int d\tau \int_0^L d\sigma \left[(\partial \vec{X})^2 + \Lambda(\vec{X}^2 - I) + \text{fermions} \right]$$

Metsaev-Tseytlin sigma model of GS superstring

CFT/AdS duality

$$\mathcal{O}(x) = \text{Tr} [D D \Psi \Psi \Phi \Phi D \Psi \dots] (x)$$



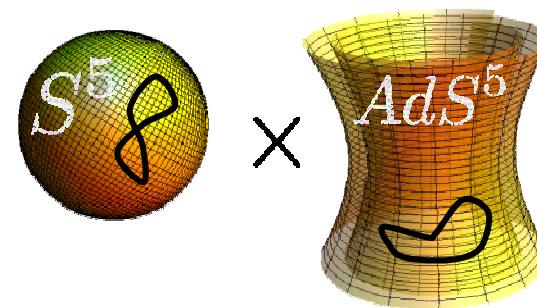
$$\mathcal{O}_A(\Lambda x) \rightarrow \Lambda^{\Delta_A(\lambda)} \mathcal{O}_A(x)$$

$$\hat{D} \mathcal{O} = \Delta \mathcal{O}$$

Dilatation operator
as a spin chain hamiltonian

Dimension of a local operator

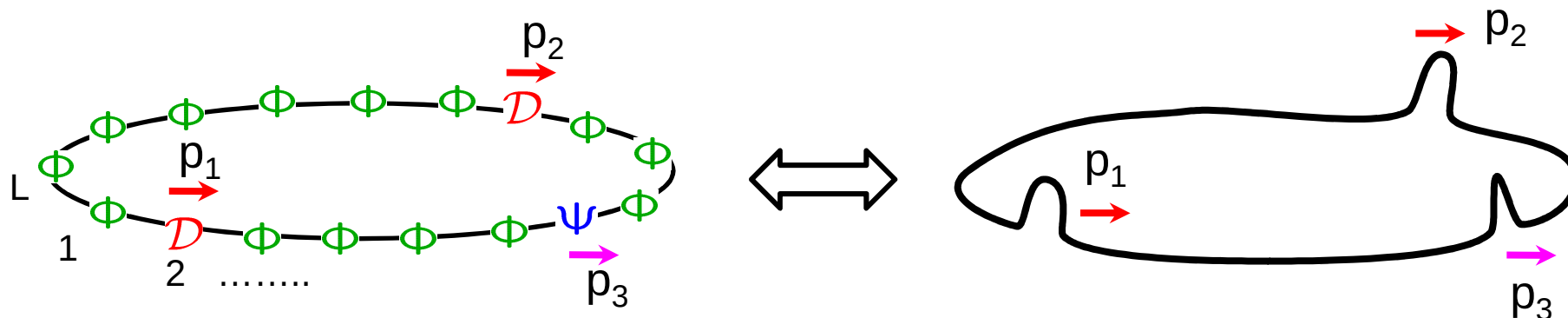
$$\lambda = g_{YM}^2 N = \frac{R^4}{\alpha'^2} = g_s N$$



Energy of a string state

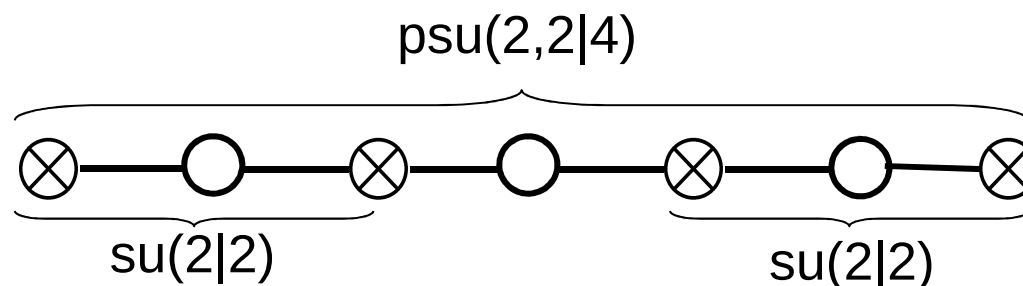
Global superconformal symmetry $\rightarrow \text{psu}(2,2|4) \leftarrow$ isometry of background*

Integrability and S-matrix



- Light cone gauge breaks the global and world-sheet Lorentz symmetries :

$$\text{psu}(2,2|4) \implies \text{su}(2|2) \times \text{su}(2|2)$$



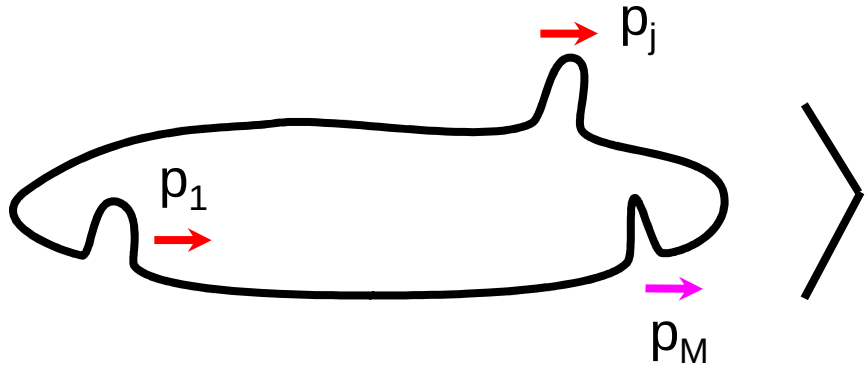
Beisert,
Staudacher
Janik
Eden,Beisert,Staudacher...

- S-matrix (Hubbard R!) of AdS/CFT via bootstrap of A.&A.I. Zamolodchikov:

$$\hat{S}_{\text{PSU}(2,2|4)}(p_1, p_2) = S_0^2(p_1, p_2) \times \hat{S}_{\text{SU}(2|2)}(p_1, p_2) \times \hat{S}_{\text{SU}(2|2)}(p_1, p_2)$$

- Asymptotic integrability : factorized scattering, asymptotique Bethe ansatz...

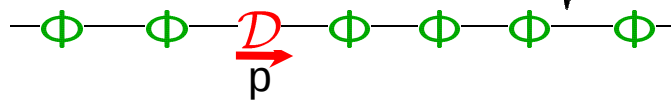
Asymptotic Bethe Ansatz (ABA)

$$0 = \left(e^{iLp_j} \prod_{k \neq j}^M \hat{S}(p_j, p_k) - 1 \right) \left| \text{Diagram} \right\rangle$$


- This periodicity condition is diagonalized by nested Bethe ansatz
- Asymptotic dispersion relation for dimensions of one particle states

Santambrogio,Zanon'02
Beisert,Dippel,Staudacher'04

$$E = \Delta - L = \sum_{j=1}^M \epsilon(p_j) + O(e^{-cL})$$

$$\epsilon(p) = \sqrt{1 + \lambda \sin^2 \frac{p}{2}}$$


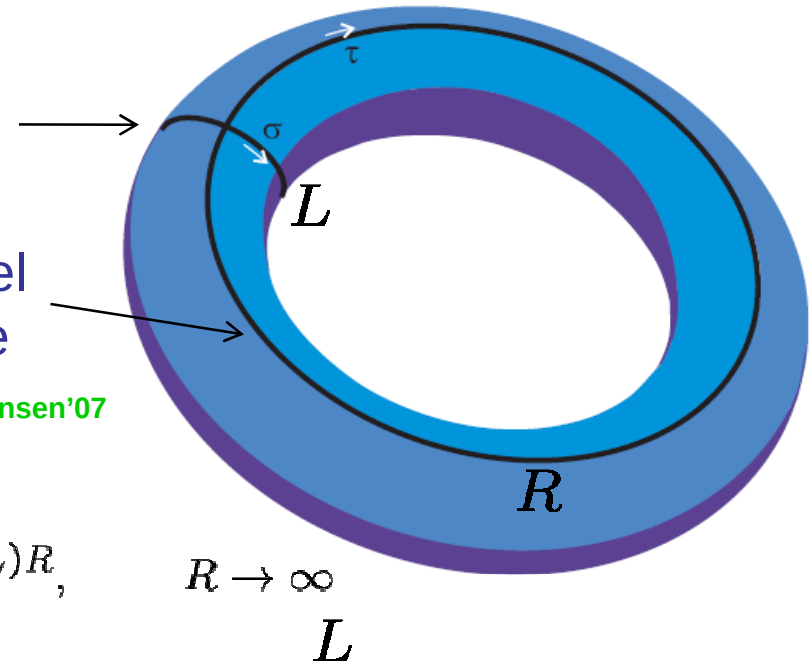
finite size corrections
Important for short operators!

Zamolodchikov TBA for short operators

original model
on L circle

"mirror" model
on R circle

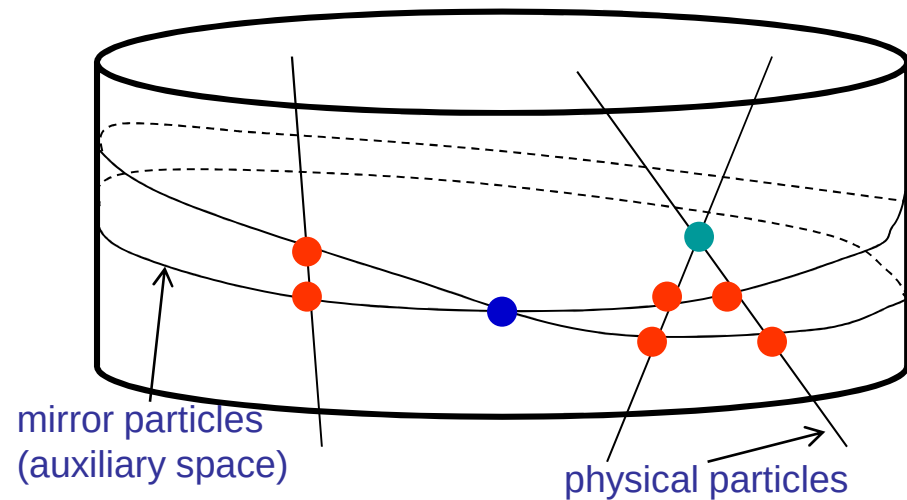
Ambjorn, Janik, Chrisijansen'07
Arutyunov, Frolov'08



$$\mathcal{Z} = \sum_n e^{-E_n^{mir}(R)L} = \sum_n e^{-E_n^{ph}(L)R} \rightarrow e^{-E_0^{ph}(L)R}, \quad R \rightarrow \infty, \quad L$$

- 2d vacuum consists of a dense gas of virtual (mirror) "particles" and a few physical particles.
- Scattering: mir-mir, mir-phys, phys-phys
- Fusion is simple on a u-plane (rapidity)
- Large R : mirror rapidities localize on poles of S-matrix $\rightarrow$ bound states

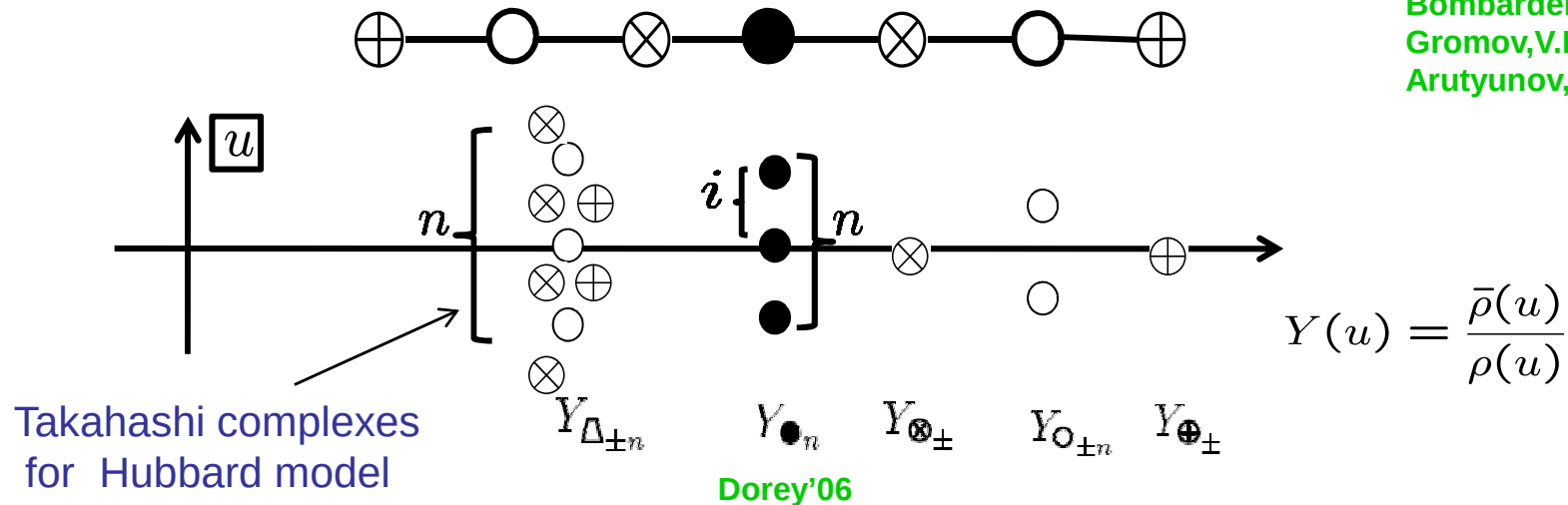
$$S_{ab}^{cd}(p(u), p(v)) \sim \frac{1}{u - v + i}$$



reminds the spin chain transfer matrix

Bound states and TBA in AdS/CFT

Gromov,V.K.,Vieira'09
Bombardelli,Fioravanti,Tateo'09
Gromov,V.K.,Kozak,Vieira'09
Arutyunov,Frolov'09



- Dispersion in Zhukovski parametrization:

$$\epsilon(u) = \sqrt{\lambda} \left(\frac{i}{x^-(u)} - \frac{i}{x^+(u)} \right), \quad p(u) = i \log \left(\frac{x^-(u)}{x^+(u)} \right)$$

- Mirror (Wick) transformation :

$$\epsilon \rightarrow i\tilde{p} \equiv \epsilon^*, \quad p \rightarrow i\tilde{\epsilon} \equiv p^*$$

$$\epsilon = \sqrt{1 + 8\lambda \sin^2 \frac{p}{2}} \rightarrow \sqrt{8\lambda} \sinh \frac{\tilde{\epsilon}}{2} = \sqrt{1 + \tilde{p}^2}$$

$$x^+(u) \rightarrow x^+(u), \quad x^-(u) \rightarrow x^-(u^*) \sim \frac{1}{x^-(u)}$$

Zhukovski map:

$$x + \frac{1}{x} = \frac{u}{\sqrt{\lambda}}$$

$$x^{\pm} + \frac{1}{x^{\pm}} = \frac{u \pm i/2}{\sqrt{\lambda}}$$

Y-system for the spectrum of AdS/CFT

- TBA: Minimizing free energy at finite temperature $T=1/L$: TBA eqs. and Y-system

$$f = \sum_{(a,s) \in \text{T-hook}} \int \left[\delta_{s,0} L \rho_{a,0} \tilde{p}_a - \rho_{a,s} \log \left(1 + \frac{\bar{\rho}_{a,s}}{\rho_{a,s}} \right) - \bar{\rho}_{a,s} \log \left(1 + \frac{\rho_{a,s}}{\bar{\rho}_{a,s}} \right) \right]$$

- Extremum condition where $\rho_{a,s} + \bar{\rho}_{a,s} = \delta_{s,0} \tilde{\epsilon}'_a + K_{a',s';a,s} * \rho_{a',s'}$

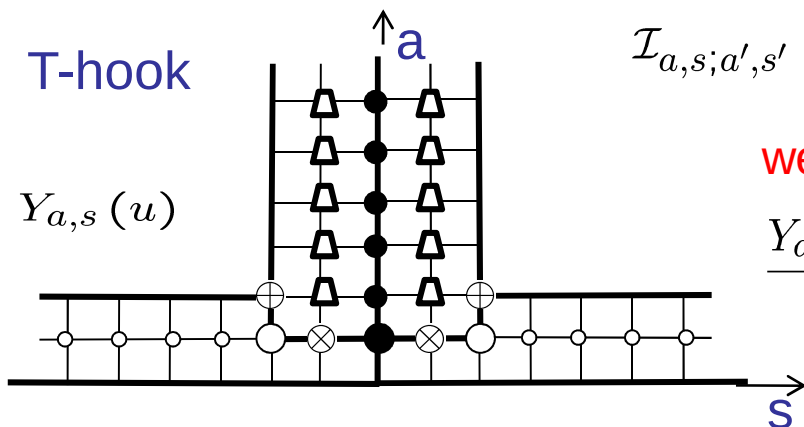
$$\log Y_{a,s} = \delta_{s,0} L \tilde{p}'_a + K_{a,s;a',s'} * \log \left(1 + Y_{a',s'}(u) \right) \quad \text{only ratios enter!}$$

$$Y_{a,s}(u) = \frac{\bar{\rho}_{a,s}(u)}{\rho_{a,s}(u)}$$

- Act by discrete Laplace operator on both parts of this eq.

$$\Delta f_a(u) \equiv f_a(u + i/2) + f_a(u - i/2) - f_{a+1}(u) - f_{a-1}(u)$$

- Using $\Delta K(u, v) = \delta(u - v) \mathcal{I}$ where $\mathcal{I}$ is the incidence matrix of T-hook

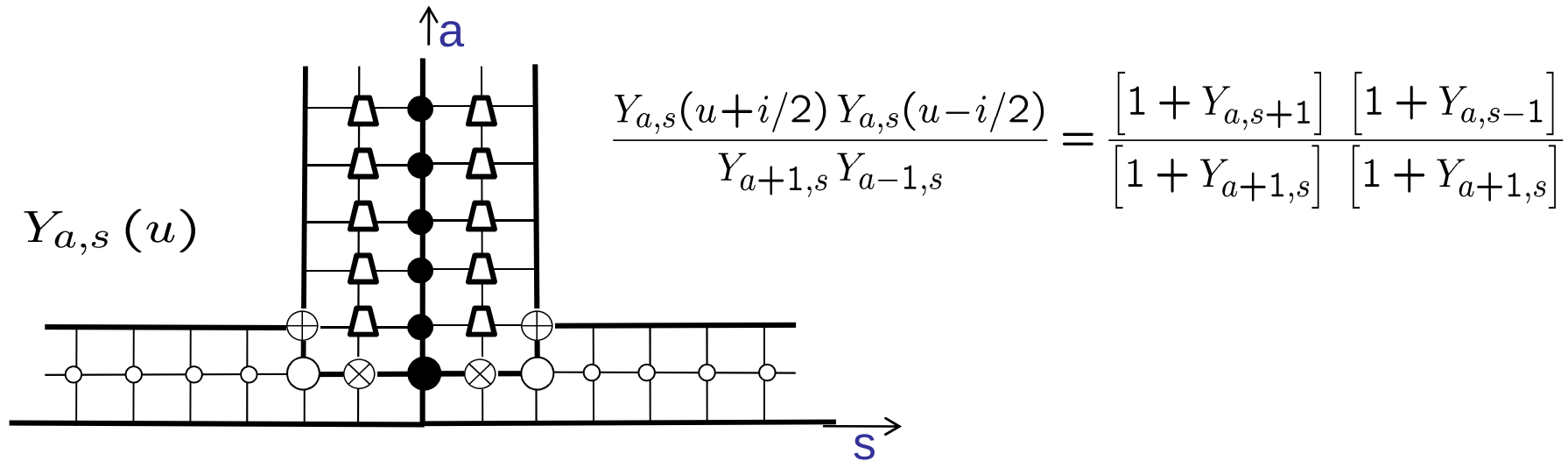


$$\mathcal{I}_{a,s;a',s'} = \delta_{a,a'} (\delta_{s',s+1} + \delta_{s',s-1}) - \delta_{s,s'} (\delta_{a',a+1} + \delta_{a',a-1})$$

we get the universal Y-system!

$$\frac{Y_{a,s}(u+i/2) Y_{a,s}(u-i/2)}{Y_{a+1,s} Y_{a-1,s}} = \frac{[1 + Y_{a,s+1}][1 + Y_{a,s-1}]}{[1 + Y_{a+1,s}][1 + Y_{a-1,s}]}$$

Y-system for excited states of AdS/CFT



- Anomalous dimensions: $\Delta = f_{min} = \int \frac{du}{2\pi i} \partial_u \epsilon_a(u) \log(1 + Y_{a,0}(u))$
- u_j obey the exact Bethe eqs: $Y_{1,0}(u_j) + 1 = 0, \quad j = 1, 2, \dots, M$

Bazhanov, Lukyanov, Zamolodchikov
Dorey, Tateo
Fioravanti, Quattrini, Ravanini

For AdS/CFT Gromov, V.K., Kozak, Vieira

Y-system and Hirota eq.: discrete integrable dynamics

- Relation of Y-system to T-system (Hirota equation) $Y_{a,s} = \frac{T_{a,s+1} T_{a,s-1}}{T_{a+1,s} T_{a-1,s}}$
 . (the Master Equation of Integrability!)

$$T_{a,s}(u+i/2)T_{a,s}(u-i/2) = T_{a,s-1}(u) T_{a,s+1}(u) + T_{a+1,s}(u) T_{a-1,s}(u)$$

- Discrete classical integrable dynamics!
- General solution of Hirota equations with Dirichlet b.c. on a line:

$$T_{a,s}(u) = \text{Det}_{1 \leq k,j \leq a} T_{1,s+k-j}(u+k+j)$$

- General solution in a strip of width N:

$$T_{a,s}(u) = \left| \begin{array}{cccc} \bar{q}_1(u-s-1) & \bar{q}_2(u-s-1) & \cdots & \bar{q}_N(u-s-1) \\ \cdots & \cdots & \cdots & \cdots \\ \bar{q}_1(u-s-a) & \bar{q}_2(u-s-a) & \cdots & \bar{q}_N(u-s-a) \\ q_1(u+s-a-1) & q_2(u+s-a-1) & \cdots & q_N(u+s-a-1) \\ \cdots & \cdots & \cdots & \cdots \\ q_1(u+s-N) & q_2(u+s-N) & \cdots & q_N(u+s-N) \end{array} \right| \begin{array}{l} \left. \vphantom{\begin{array}{c} \bar{q}_1(u-s-1) \\ \cdots \\ \bar{q}_1(u-s-a) \end{array}} \right\} a \\ \left. \vphantom{\begin{array}{c} q_1(u+s-a-1) \\ \cdots \\ q_1(u+s-N) \end{array}} \right\} N-a \end{array}$$

Krichever, Lipan,
Zabrodin, Wiegmann

- Gauge transformations of T...

$$T_{a,s}(u) \rightarrow G_1\left(u+i\frac{a+s}{2}\right) G_2\left(u+i\frac{a-s}{2}\right) G_3\left(u-i\frac{a+s}{2}\right) G_4\left(u-i\frac{a-s}{2}\right) T_{a,s}(u)$$

Hirota equation and characters of supergroups

- Jacobi-Trudi formula for $SL(K|M)$ characters general irrep $\lambda = \{\lambda_1, \lambda_2, \dots, \lambda_a\}$

$$T_{\{\lambda\}}[g] = \det_{1 \leq i, j \leq a} T_{1, \lambda_i - i + j}[g], \quad g = \text{diag}\{x_1, \dots, x_K | y_1, \dots, y_M\} \in SL(K|M).$$

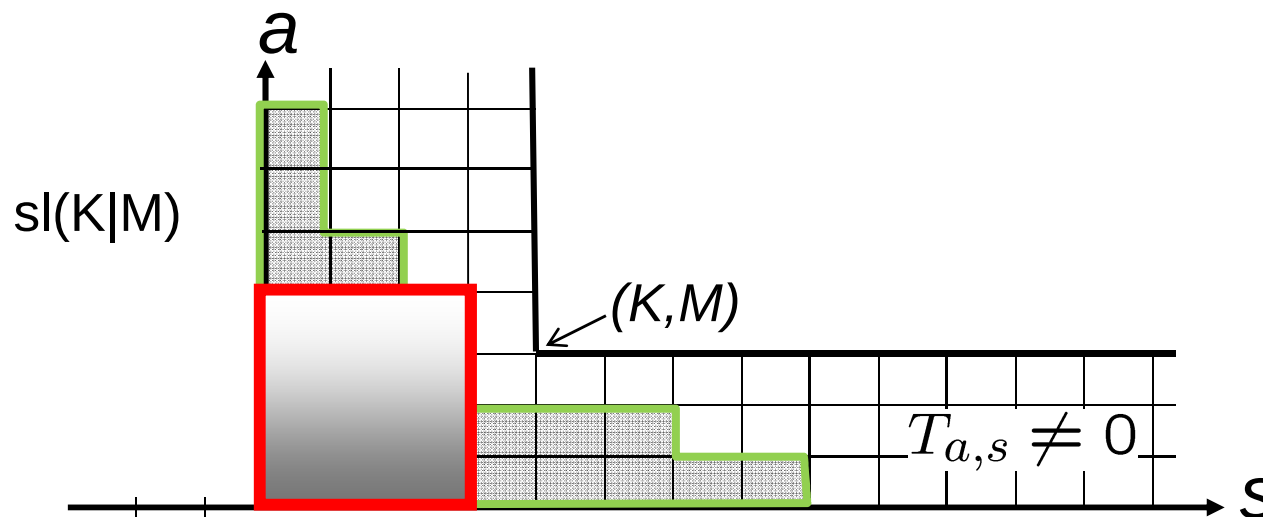
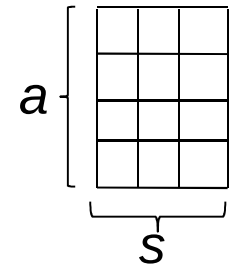
$T_{1,s}[g]$ are super-characters in symmetric irreps generated by



$$w(z) \equiv \text{sdet} (1 - zg)^{-1} = \sum_{s=1}^{\infty} T_{1,s}(g) z^s$$

- For rectangular Young tableaux (a,s) Hirota eq. with **fat hook** b.c.:

$$T_{a,s}^2 = T_{a+1,s} T_{a-1,s} + T_{a,s+1} T_{a,s-1}$$



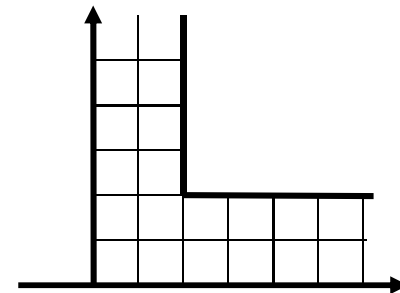
Generalization
to super-spin chains:

V.K., Zabrodin, Sorin'07
Tsuboi'97

Supergroups: Fat Hook of $SU(4|4)$ and T-hook of $PSU(2,2|4)$

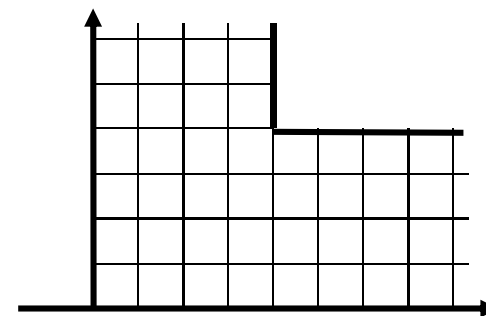
$SU(2|2)$

$$w_{SU(2|2)}(z; g) = \frac{(1 - zy_1)(1 - zy_2)}{(1 - zx_1)(1 - zx_2)} = \sum_{j=0}^{\infty} T_{1,s} z^s$$



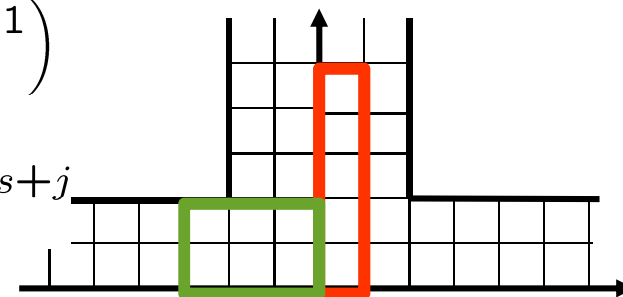
$SU(4|4)$

$$\begin{aligned} w_{SU(4|4)}(z; g \otimes g') &= w_{SU(2|2)}(z; g) \times w_{SU(2|2)}(z; g') \\ &= \sum_{s=0}^{\infty} z^s \sum_{j=0}^s T_{1,j} T'_{1,s-j} \end{aligned}$$



$SU(2,2|4)$

$$\begin{aligned} w_{SU(2,2|4)}(z; g \otimes g') &= w_{SU(2|2)}(z; g) \times w_{SU(2|2)}\left(\frac{1}{z}; g'^{-1}\right) \\ &= \frac{y_3 y_4}{x_3 x_4} \sum_{s=-\infty}^{\infty} z^s \sum_{j=\text{Min}(-s, 0)}^{\infty} \bar{T}'_{1,j} T_{1,s+j} \end{aligned}$$



- Using these characters one solves AdS/CFT quasiclassically
- Generalization to full quantum case (u-dependence) possible!

Gromov
V.K.
Tsuboi

Unitary highest weight representations of T-hook of non-compact superalgebra $u(M_1, M_2 | N)$

Generators:

$$[E_{ab}, E_{cd}] = \delta_{bc} E_{ad} - (-1)^{(p_a + p_b)(p_c + p_d)} \delta_{da} E_{cb},$$

Dual basis : $\varepsilon_i(E_{jj}) = \delta_{ij}$

$$(\varepsilon_i | \varepsilon_j) = p_i \delta_{ij}, \quad p_i = (+1_{M_1}, -1_N, +1_{M_2})$$

Simple roots:

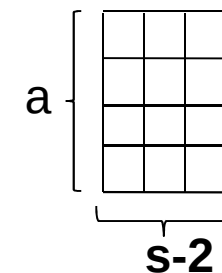
$$\alpha_i = \varepsilon_i - \varepsilon_{i+1}, \quad i = 1, 2, \dots, N + M - 1$$

Cartan matrix :

$$C_{ij} = (\alpha_i, \alpha_j)$$

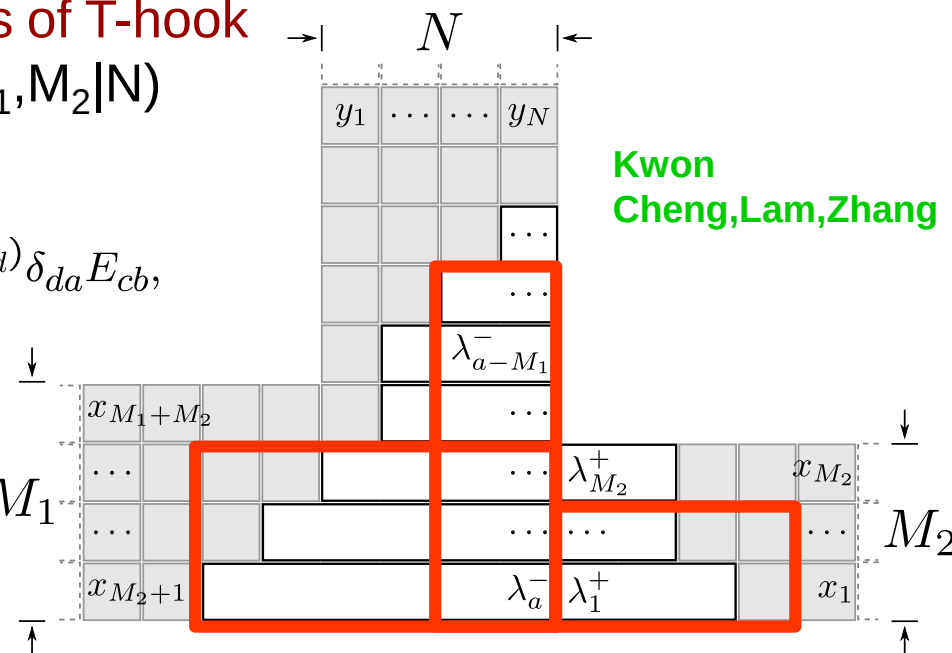
• Highest weight Λ for $u(2, |4|, 2)$

$$\Lambda = \begin{cases} a(-\varepsilon_1 - \varepsilon_2) + (s+2) \sum_{i=7}^{a+6} \varepsilon_i, & s < -2, \quad 0 \leq a \leq 2 \\ a(-\varepsilon_1 - \varepsilon_2 + \sum_{i=3}^{s+4} \varepsilon_i), & -2 \leq s \leq 2, \quad 0 \leq a \\ a(-\varepsilon_1 - \varepsilon_2 + \sum_{i=3}^6 \varepsilon_i) + (s-2) \sum_{i=7}^{a+6} \varepsilon_i, & 2 < s, \quad 0 \leq a \leq 2 \end{cases}$$



• Unitarity : $\eta E_{ab}^\dagger \eta = E_{ba}, \quad \eta = \text{diag}(-1_{M_1}, +1_N, +1_{M_2})$

*



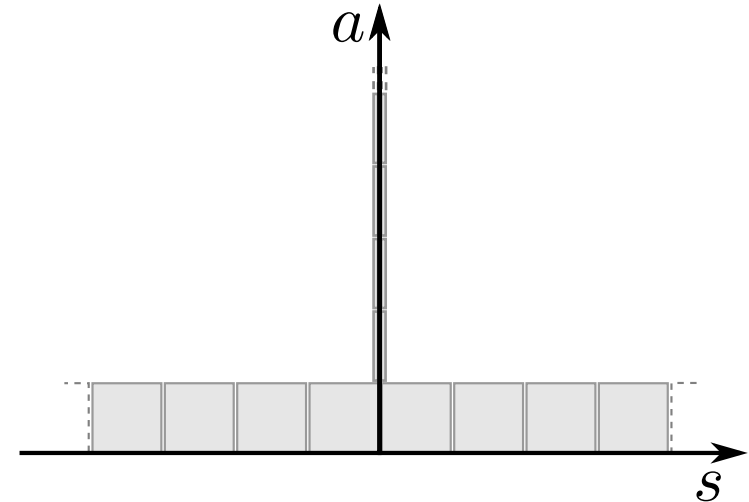
Kwon
Cheng, Lam, Zhang

Unitary non-compact representations of T-hook of $u(1,1)$

$$[h_3, h_{\pm}] = \pm 2h_{\pm} \quad , \quad [h_-, h_+] = h_3$$

highest
weight

$$\Lambda = \begin{cases} -\varepsilon_1 + |s|\varepsilon_2 & , \quad s > 0 \quad , \quad a = 1 \\ -\varepsilon_1 - |s|\varepsilon_1 & , \quad s < 0 \quad , \quad a = 1 \\ -a\varepsilon_1 & , \quad s = 0 \quad , \quad a > 0 \end{cases}$$



$$h_+ = \begin{pmatrix} 0 & a_0 & 0 & 0 & \cdots \\ 0 & 0 & a_1 & 0 & \cdots \\ 0 & 0 & 0 & a_2 & \cdots \\ 0 & 0 & 0 & 0 & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}, \quad h_- = \begin{pmatrix} 0 & 0 & 0 & 0 & \cdots \\ -a_0 & 0 & 0 & 0 & \cdots \\ 0 & -a_1 & 0 & 0 & \cdots \\ 0 & 0 & -a_2 & 0 & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}, \quad h_3 = \begin{pmatrix} -|s| - 1 & 0 & 0 & \cdots \\ 0 & -|s| - 3 & 0 & \cdots \\ 0 & 0 & -|s| - 5 & \cdots \\ \cdots & \cdots & \cdots & \cdots \end{pmatrix}$$

$$a_j = i\sqrt{(j+1)(j+|s|+1)}$$

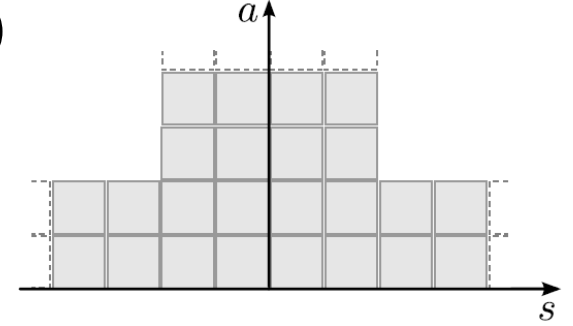
• Solution of Hirota for $U(1,1)$ characters

$$T_{1,s} \equiv T_s^{(1+1)} = \frac{1}{x_2} \sum_{j=\max(0,-s)}^{\infty} T_{s+j}^{(1)}(x_1) T_j^{(1)}(1/x_2) = \sum_{j=\max(0,-s)}^{\infty} x_1^{s+j} x_2^{-j-1}$$

$$T_{1,s} \equiv T_s^{(1+1)}(x_1, x_2) = \begin{cases} \frac{x_1^s}{x_2 - x_1} & , \quad s > 0 \\ \frac{x_2^s}{x_2 - x_1} & , \quad s \leq 0 \end{cases} \quad T_{a,0} = \frac{x_2^{1-a}}{x_2 - x_1}$$

*

Character solution of T-hook for $u(2,2|4)$



- Solution in finite determinants
(analogue of the 1-st Weyl formula)

$$T_{a,s} = \begin{cases} (-1)^{(a+1)s} \left(\frac{x_3 x_4}{y_1 y_2 y_3 y_4} \right)^{s-a} \frac{\det \left(S_i^{\theta_{j,s+2}} y_i^{j-4-(a+2)\theta_{j,s+2}} \right)_{1 \leq i,j \leq 4}}{\det \left(S_i^{\theta_{j,0+2}} y_i^{j-4-(0+2)\theta_{j,0+2}} \right)_{1 \leq i,j \leq 4}}, & a \geq |s| \\ \frac{\det \left(Z_i^{(1-\theta_{j,a})} x_i^{2-j+(s-2)(1-\theta_{j,a})} \right)_{1 \leq i,j \leq 2}}{\det \left(Z_i^{(1-\theta_{j,0})} x_i^{2-j+(0-2)(1-\theta_{j,0})} \right)_{1 \leq i,j \leq 2}}, & s \geq +a \end{cases}$$

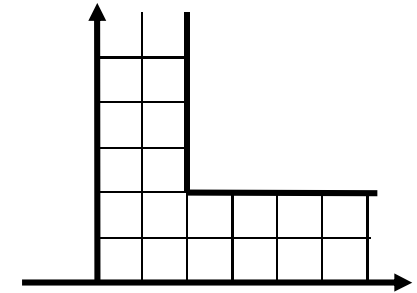
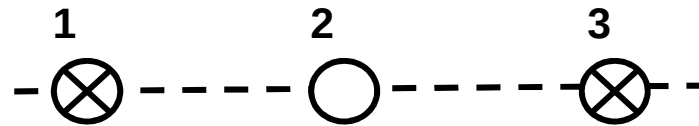
$$\theta_{j,s} = \begin{cases} 1, & j > s \\ 0, & j \leq s \end{cases}$$

$$S_i = \frac{(y_i - x_3)(y_i - x_4)}{(y_i - x_1)(y_i - x_2)}$$

$$Z_i = \frac{(x_i - y_1)(x_i - y_2)(x_i - y_3)(x_i - y_4)}{(x_i - x_3)(x_i - x_4)},$$

Solving full quantum Hirota: SU(2|2) example

- Generating functional of SU(2|2) irreps:



$$\mathcal{W} = [1 - \mathbf{T}^1 D] [1 - \mathbf{T}^2 D]^{-1} [1 - \mathbf{T}^3 D]^{-1} [1 - \mathbf{T}^4 D]_{(1)} , \quad D = e^{-i\partial_u}$$

$$\mathcal{W}_R = \sum_{s=0}^{\infty} \mathbf{T}_{1,s}^R \left(u + i \frac{1-s}{2} \right) D^s , \quad \mathcal{W}_R^{-1} = \sum_{a=0}^{\infty} (-1)^a \mathbf{T}_{a,1}^R \left(u + i \frac{1-a}{2} \right) D^a .$$

- T's are parametrized through 4(+1) functions

$$\mathbf{T}^1(u) = h \frac{Q_1^- B_4^+}{Q_1^+ B_4^-},$$

$$\mathbf{T}^2(u) = h \frac{Q_1^- Q_2^{++}}{Q_1^+ Q_2^-},$$

$$\mathbf{T}^3(u) = h \frac{Q_2^{--} Q_3^+}{Q_2^- Q_3^-},$$

$$\mathbf{T}^4(u) = h \frac{Q_3^+ R_4^-}{Q_3^- R_4^+}$$

Back to asymptotic Bethe eqs. and S-matrix: $L \rightarrow \infty$

Gromov, V.K., Vieira'09

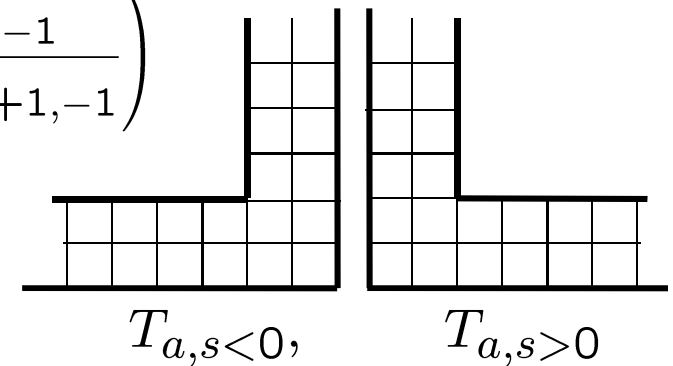
- From TBA: $Y_{a \geq 1, 0} \sim \left(\frac{x^{[-a]}}{x^{[+a]}} \right)^L \rightarrow 0$

- It is a spin chain limit:

$$\frac{Y_{a,0}^+ Y_{a,0}^-}{Y_{a-1,0} Y_{a+1,0}} \simeq \left(\frac{T_{a,1}^+ T_{a,1}^-}{T_{a-1,1} T_{a+1,1}} \right) \left(\frac{T_{a,-1}^+ T_{a,-1}^-}{T_{a-1,-1} T_{a+1,-1}} \right)$$

$$1 + Y_{a,s} = \frac{T_{a,s}^+ T_{a,s}^-}{T_{a+1,s} T_{a-1,s}}$$

- T-system splits into two $SU(2|2)_{L,R}$ wings:



- Solving this discrete Laplace eq. in (a,u) -variables we get

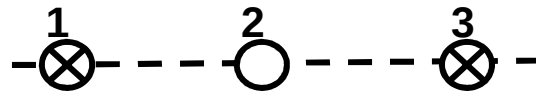
$$Y_{a,0}(u) \simeq \underbrace{\left(\frac{x^{[-a]}}{x^{[+a]}} \right)^L \frac{\phi^{[-a]}}{\phi^{[+a]}}}_{\text{zero mode (dressing factor...)}} T_{a,-1}^L T_{a,1}^R$$

transfer matrices of $SU(2|2)$
(fixed from analyticity)

Asymptotic Bethe Ansatz $L \rightarrow \infty$: concise form

- Fundamental transfer matrix for $SU(2|2)_{L,R}$

$$\mathbf{T}_{L,R}(u) = -\mathbf{T}_{L,R}^1(u) + \mathbf{T}_{L,R}^2(u) + \mathbf{T}_{L,R}^3(u) - \mathbf{T}_{L,R}^4(u)$$



- Scalar dressing factor and T-functions:

1. Analyticity (each T has ≤ 2 cuts)

2. Y's are real functions

$$\frac{\phi^-}{\phi^+} = \frac{Q_4^{--}}{Q_4^{++}} \cdot \prod_{j=1}^{K_4} \left[\frac{\frac{1}{x^+} - x_{4,j}^-}{\frac{1}{x^-} - x_{4,j}^+} \right] \cdot S_0^2$$

$$S_0(u) = \prod_{j=1}^{K_4} \sigma(x(u), x_{4,j})$$

- Crossing equation (from reality of Y's)

$$\sigma(x, y) \cdot \sigma(x^*, y) = \frac{y^- x^- - y^+}{y^+ x^+ - y^-} \cdot \frac{1/x^- - y^+}{1/x^+ - y^-}$$

$$\begin{aligned} \mathbf{T}^1(u) &= h \frac{Q_1^-}{Q_1^+} \prod_{j=1}^{K_4} \frac{\frac{1}{x^+} - x_{4,j}^-}{\frac{1}{x^+} - x_{4,j}^+}, \\ \mathbf{T}^2(u) &= h \frac{Q_1^- Q_2^{++}}{Q_1^+ Q_2^+}, \\ \mathbf{T}^3(u) &= h \frac{Q_2^- Q_3^+}{Q_2^+ Q_3^-}, \\ \mathbf{T}^4(u) &= h \frac{Q_3^+}{Q_3^-} \prod_{j=1}^{K_4} \frac{x^- - x_{4,j}^+}{x^- - x_{4,j}^-} \end{aligned}$$

Defs:

$$Q_l(u) = \prod_{j=1}^{K_l} (u - u_{l,j})$$

$$l = 1, 2, 3(5, 6, 7)$$

$$h = \prod_{j=1}^{K_4} \frac{x^- - x_{4,j}^-}{x^- - x_{4,j}^+} \prod_{j=1}^{K_1} \frac{\frac{1}{x^+} - x_{1,j}}{\frac{1}{x^-} - x_{1,j}} \prod_{j=1}^{K_3} \frac{\frac{1}{x^-} - x_{3,j}}{\frac{1}{x^+} - x_{3,j}}$$

Janik

Beisert, Eden, Staudacher,

Volin

- Beisert-Staudacher ABA – from $Y_{1,0}(u_j) + 1 = 0$
- Auxiliary Bethe eqs. from regularity of $\mathbf{T}_{L,R}(u)$

Konishi at 4 loops (one wrapping) from Y-system

- We have to calculate $Y_{a,0}$ in mirror for $L=2$
- Dressing factor does not contribute at this order
- Two physical roots defined from central ABA eq.

$$\left(\frac{x^-(u_{\pm})}{x^+(u_{\pm})}\right)^2 \frac{u_{\pm} - i/2}{u_{\pm} + i/2} = 1 \Rightarrow u_{\pm} = \pm \frac{1}{2\sqrt{3}}$$

- Mirror dispersion:

$$\epsilon^{mir}(u) = -\log(x^{[-a]}x^{[+a]}) = \log \frac{g^4}{(u^2 + a^2/4)^2} + \mathcal{O}(g^5)$$

$$Y_{a,0}^{mir}(u) = g^8 \left(3 \cdot 2^7 \frac{3a^2 + 12au^2 - 4a}{(a^2 + 4u^2)^2} \right)^2 \frac{1}{y_a(u)y_{-a}(u)}$$

$$y_a(u) = 9a^2 - 36a^3 + 72u^2a^2 + 60a^2 - 144u^2a - 48a + 144u^4 + 48u + 16$$

$$(\Delta_K - 2)_{wrap} \simeq \sum_{a=1}^{\infty} \int_{-\infty}^{\infty} du \epsilon_a^{mir}(u) \cdot Y_{a,0}^{mir}(u) =$$

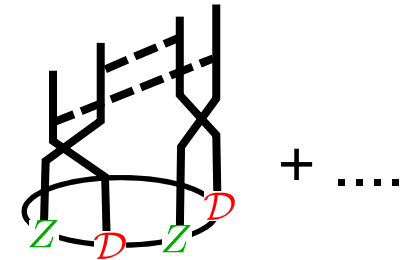
$$g^8(324 + 864\zeta(3) - 1440\zeta_5) + \mathcal{O}(g^{10})$$

~131.000 4D Feynman graphs!

- Now 5 loops available:

$$\Delta_{Konishi} = 2 + 12\lambda - 48\lambda^2 + 336\lambda^3 - [(2820 + 288\zeta(3)) + (324 + 864\zeta(3) - 1440\zeta(5))] \lambda^4 \\ + [(26508 + 4320\zeta(3) + 2880\zeta(5)) + (-11340 + 2592\zeta(3) - 5184\zeta(3)^2 - 11520\zeta(5) + 30240\zeta(7))] \lambda^5 + \mathcal{O}(\lambda^6)$$

$$\mathcal{O}_{Konishi} = \text{Tr} [\mathcal{D}, Z]^2$$



Minahan, Zarembo,

Beisert, Kristijanssen,
Staudacher

Beisert, Dippel, Staudacher
Beisert, Eden, Staudacher,

Bajnok, Janik, Lukowski

Bern, Dixon, Kosover, Smirnov
Fiamberti, Zieg, Zanon,
Santambogio
Velizhanin

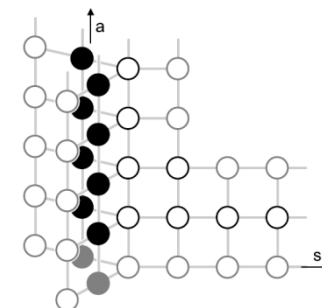
Gromov, V.K., Vieira
Arutyunov, Frolov
Balog, Hegedus

Conclusions and problems

Done:

- Y-système and TBA for exact spectrum of AdS/CFT
- Konishi dimension at “all” couplings (numerics)
- Tested at weak (4 loops) and strong (one loop) coupling
- Two loops at strong coupling predicted
- 5-loops done
- Y-system for CP3xAdS4 available and checked

$$\Delta_{Konishi} = 2\lambda^{1/4} + 2\lambda^{-1/4} + O(\lambda^{-3/4})$$



Gromov,V.K.,Vieira
Minahan,OhlssonSax,Sieg
Bajnok, Janik
Bombardelli,Fioravanti,Tateo
Gromov,Levkovich-Maslyuk

To do:

- Derivation of Y-system from SYM?
It would prove both integrability and AdS/CFT correspondance
- Equation of Destri-DeVega-type, using Hirota discrete integrable dynamics
- Regular weak and strong coupling expansions for Konishi and other operators
- BFKL via Y-system?
- Lessons for QCD?
- Classification and solution of «all» integrable sigma models using Hirota dynamics
- New AdS/CFT's

Zarembo

End